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Polarization properties of chiral-core planar waveguides

Presenter: Warren N. Herman

Address: US Navy, NAWCAD
EO Sensors Division AIR456
Patuxent River, MD 20670

Abstract

A new form for the modal eigenvalue equations of chiral-core planar waveguides provides insight into the transition, with increasing chirality, from TE/TM (transverse electric/magnetic) modes in the achiral case to nearly right-handed and left-handed circularly polarized modes. Dramatic variation of the polarization eccentricity with thickness and frequency is discussed.



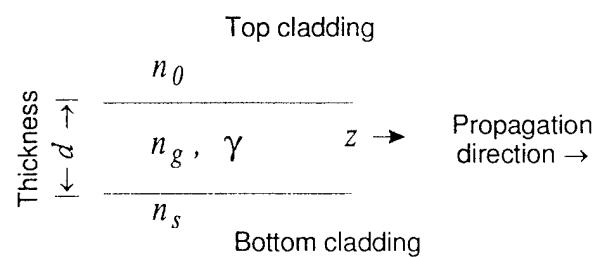
Polarization Properties of Chiral Core Planar Waveguides

W.N. Herman
U.S. Navy, NAWCAD
EO Sensors Division, Patuxent River, MD

OTF2001
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Chiral Core Optical Waveguides



ISOTROPIC CHIRAL MEDIUM

Constitutive Relations
(Drude-Born-Fedorov)

$$\vec{D} = \epsilon(\vec{E} + \gamma \nabla \times \vec{E}) \quad \vec{B} = \mu(\vec{H} + \gamma \nabla \times \vec{H})$$

Bohren's Decomposition

$$\vec{E} = \frac{1}{2}(\vec{F}^+ + \vec{F}^-) \quad \vec{H} = \frac{1}{2i}\sqrt{\frac{\epsilon}{\mu}}(\vec{F}^+ - \vec{F}^-)$$

Wave Equation

$$\nabla^2 \vec{F}^\pm + (k_0 n_\pm)^2 \vec{F}^\pm = 0$$

Eigenmodes in bulk
material circularly polarized

$$\vec{F}^\pm = \vec{E}_0 e^{i(k_0 n_\pm z - \omega t)} (\hat{x} \pm i\hat{y})$$

Refractive indices
for RH and LH
waves

$$n_\pm = \frac{n_g}{1 \pm \delta} \quad \delta = k_0 n_g \gamma \quad n_g = \sqrt{\frac{\epsilon}{\epsilon_0}} \quad k_0 = \frac{2\pi}{\lambda}$$

Rotatory Power

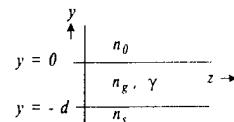
$$\rho = \frac{\pi(n_- - n_+)}{\lambda} \approx k_0 n_g \delta$$

A. Lahtakia, V.K. Varadan, and V.V. Varadan, Time-Harmonic Electromagnetic Fields in Chiral Media, Lecture Notes in Physics Series 335
(Springer-Verlag, Berlin, 1989).
I.V. Lindell, A.H. Sihvola, S.A. Tretyakov, A.J. Viitanen, Electromagnetic Waves in Chiral and Bi-isotropic Media,
(Artech House, Norwood, MA, 1994).

Modes in Chiral Asymmetric Waveguide

W. N. Herman, accepted J. Opt. Soc. Am. A

$$\vec{F}^\pm(y, z) = \vec{\Psi}^\pm(y) \exp(-ik_0 n_{\text{eff}} z)$$



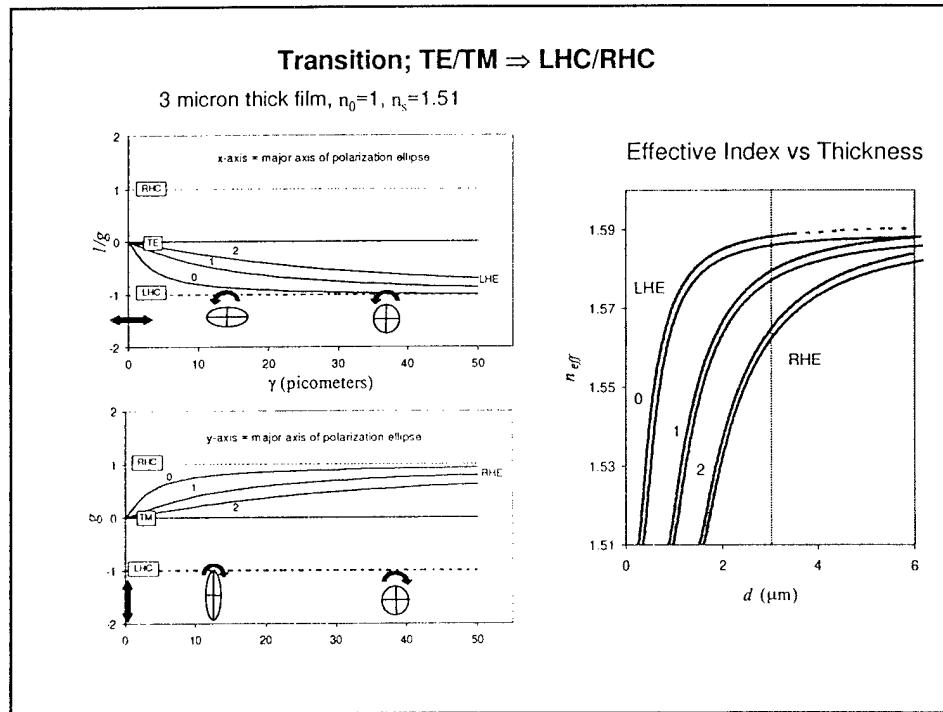
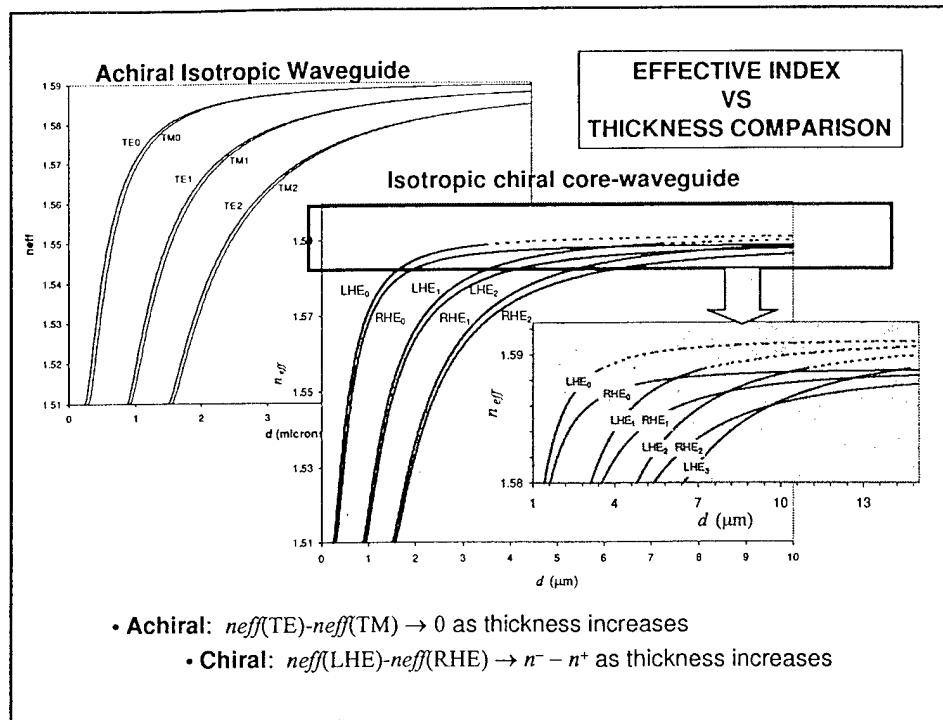
Modal equations: (3 equations to be solved simultaneously for n_{eff} , g , h)

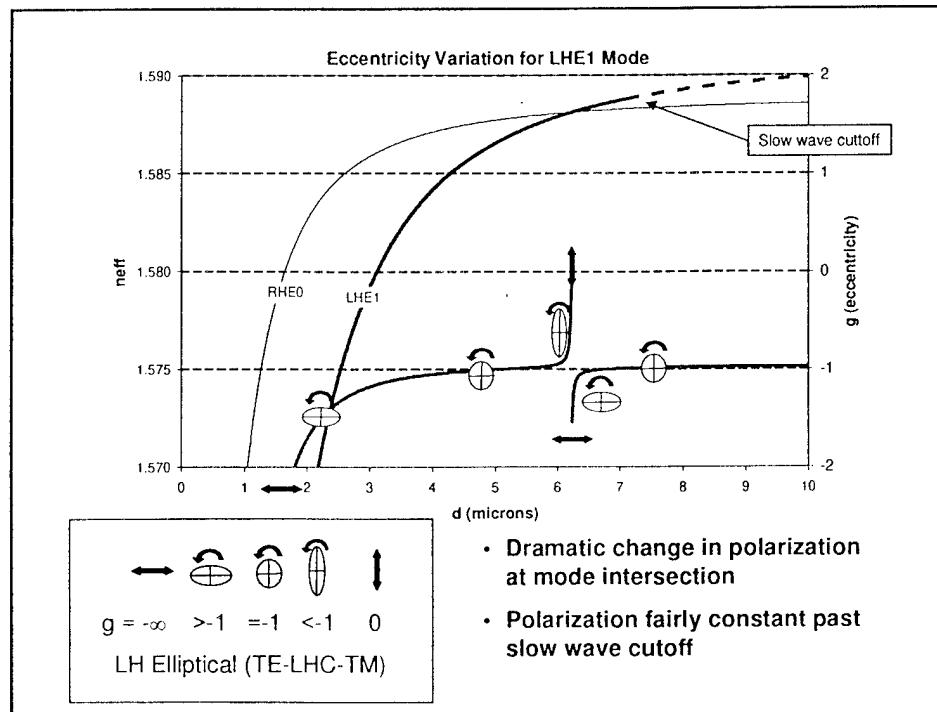
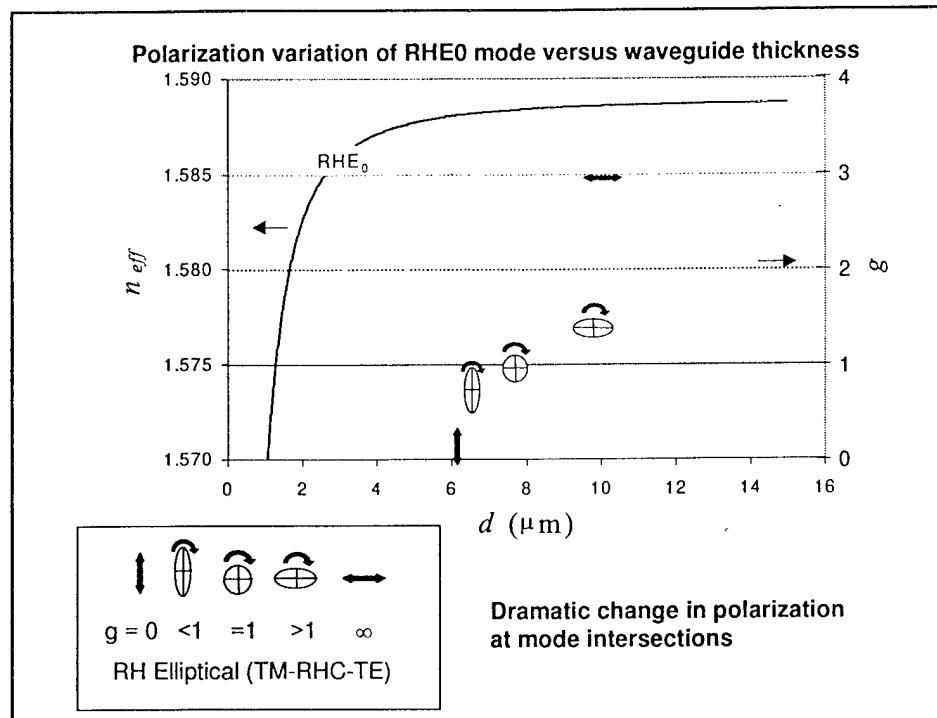
$$\begin{aligned} u^* d &= \cot^{-1} \left(\sigma_0^\pm \frac{r_0 \pm g}{1 \pm g} \right) + \cot^{-1} \left(\sigma_i^\pm \frac{r_i \pm h}{1 \pm h} \right) + m^\pm \pi \\ h(g, n_{\text{eff}}) &= \frac{(S_0^+ - S_0^-) + (S_0^+ + S_0^-)g}{(S_0^+ + S_0^-) + (S_0^+ - S_0^-)g} \end{aligned} \quad \begin{aligned} S_0^\pm &\equiv \sigma_0^\pm \sin(u^* d) - \cos(u^* d) \\ S_0^\pm &\equiv r_0 \sigma_0^\pm \sin(u^* d) - \cos(u^* d) \end{aligned}$$

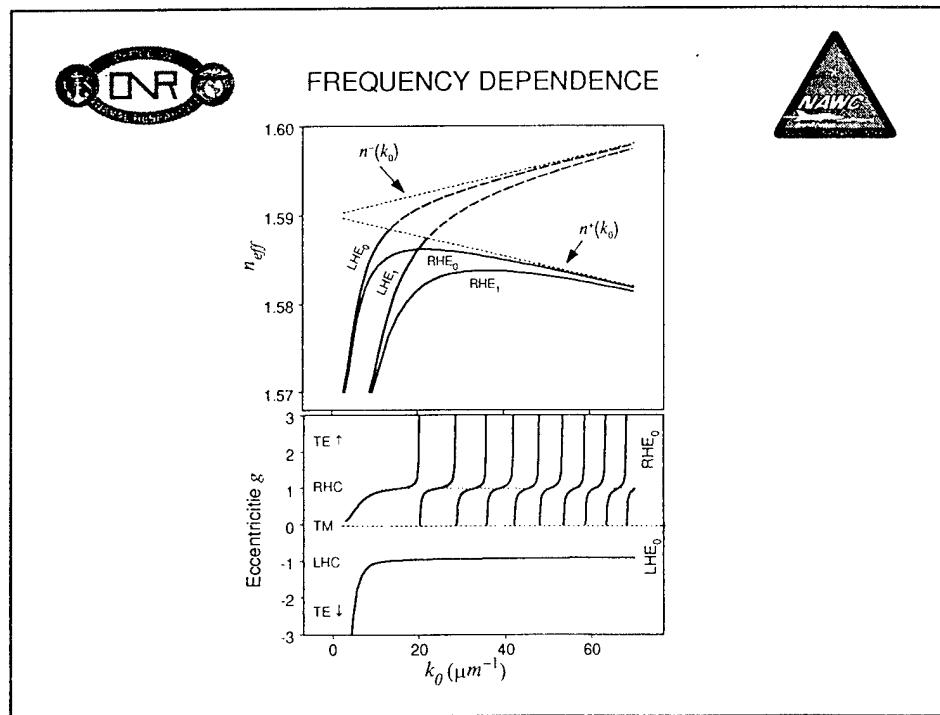
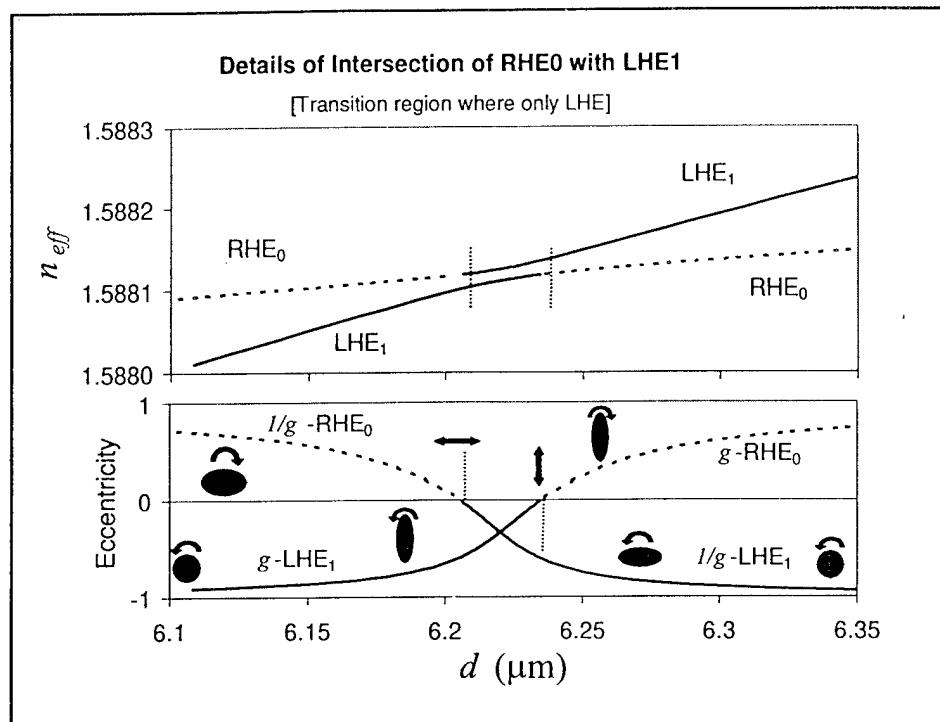
$$\begin{aligned} \sigma_0^\pm &\equiv (1 \pm \delta) \frac{u^\pm}{v}, \quad \sigma_i^\pm \equiv (1 \pm \delta) \frac{u^\pm}{w}, \quad r_0 \equiv \frac{n_0^2}{n_\pm^2}, \quad r_i \equiv \frac{n_i^2}{n_\pm^2} \\ u^\pm &\equiv k_0 \sqrt{n_\pm^2 - n_{\text{eff}}^2}, \quad v \equiv k_0 \sqrt{n_{\text{eff}}^2 - n_0^2}, \quad w \equiv k_0 \sqrt{n_{\text{eff}}^2 - n_i^2} \end{aligned} \quad n^\pm = \frac{n_g}{1 \pm \delta}$$

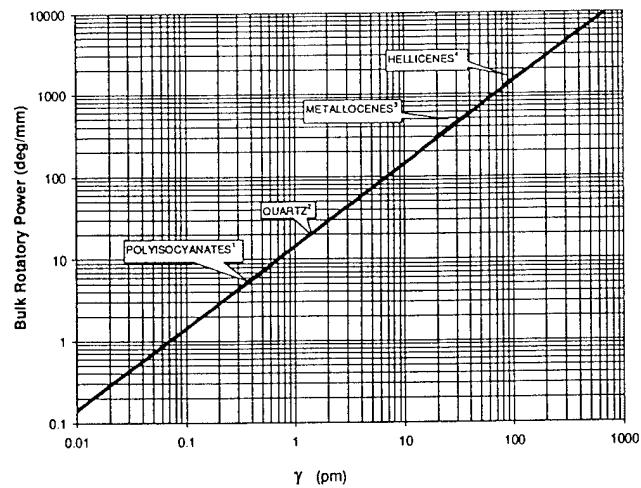
Parameters g , h determine eccentricity of polarization ellipse for transverse E-field:

$$\left. \frac{E_x}{E_z} \right|_{v \geq 0} = i \frac{n_{\text{eff}}}{n_g} \frac{1}{g} \quad \left. \frac{E_y}{E_z} \right|_{v \leq -d} = i \frac{n_{\text{eff}}}{n_g} \frac{1}{h}$$









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